

# A Test Structure to Characterize Nano-Scale Ohmic Contacts in III-V MOSFETs

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**Abstract**—We propose and demonstrate a novel test structure to characterize the electrical properties of nano-scale metal–semiconductor contacts. The structure is in essence a two-port transmission line model (TLM) with contacts in the nanometer regime. Unlike the conventional TLM, two types of Kelvin measurements are possible. When performed on devices with different contact spacing, this allows the extraction of the contact resistance, the semiconductor sheet resistance, and the metal sheet resistance. For this, a 2-D distributed resistive network model has been developed. We demonstrate this technique in Mo/n<sup>+</sup>-InGaAs contacts with contact lengths from 19 to 450 nm where we have measured an average contact resistivity of  $0.69 \pm 0.3 \Omega \cdot \mu\text{m}^2$ . For relatively long contacts ( $>110$  nm), this corresponds to an extremely small contact resistance of  $6.6 \pm 1.6 \Omega \cdot \mu\text{m}$ .

**Index Terms**—Nano contacts, TLM, contact resistivity, III-V MOSFET.

## I. INTRODUCTION

THE III-V MOSFET is one of the most promising technologies to extend Moore’s law beyond the limits of Si [1], [2]. In these devices, contact resistance remains a great challenge. At the insertion point, III-V logic transistors must have contacts in the 10–15 nm length range while delivering a total source resistance of  $50 \Omega \cdot \mu\text{m}$  [1], [3]. To meet this goal, the metal to semiconductor contact resistivity  $\rho_c$  needs to be lower than  $0.5 \Omega \cdot \mu\text{m}^2$ . Low values of  $\rho_c$  for contacts on n<sup>+</sup>-InGaAs between 0.4 and  $3.2 \Omega \cdot \mu\text{m}^2$  have been reported [4]–[7]. Currently, the contact resistance in lateral devices is mostly extracted using a TLM structure [8] or a four-terminal Kelvin test structure [9]. In a TLM, the contacts are very large, they are separated by a distance in the scale of microns and the metal resistance is neglected. However, when dealing with nano-scale contacts for future CMOS, we need a more accurate test structure capable of extracting extremely small values of contact resistance on very small contacts where the metal resistance can be significant. Addressing this need is the goal of this letter.

In this letter, we propose a new two-port TLM structure with nano-scale contacts separated by nanometer-range distances.

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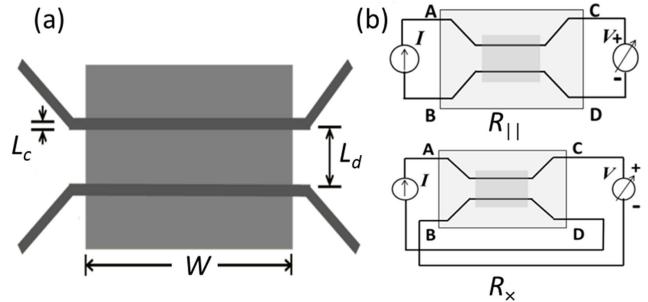


Fig. 1. (a) Schematic top-down view of a single nano-TLM test structure. (b) Two Kelvin measurement schemes used to characterize the nano-TLM test structure.

We model the structure through a 2D distributed circuit network. The model allows the simultaneous extraction of the contact resistance,  $R_c$ , the semiconductor sheet resistance,  $R_{sh}$ , and the metal sheet resistance,  $R_{shm}$ . The proposed approach is demonstrated by fabricating Mo nano-contacts on n<sup>+</sup>-InGaAs heterostructures where we have extracted an average contact resistivity as low as  $0.69 \Omega \cdot \mu\text{m}^2$ . This matches the state-of-the-art in this metal/semiconductor system [4]–[7]. The proposed test structure and model is not limited to III-V MOSFETs and can be used to analyze nano-scale contacts in any material system.

## II. NANO-TLM DESIGN AND MODELING

Fig. 1(a) shows a schematic of our proposed nano-TLM test structure. The shaded area indicates the active semiconductor region, on which there are two parallel thin metal contacts of length  $L_c$ , width  $W$ , separated by a distance  $L_d$ . A unique aspect of this structure is that the two nano-scale metal lines can be contacted at both ends separately. This enables the two Kelvin measurement schemes illustrated in Fig. 1(b).  $R_{||}$ , the parallel terminal resistance, is typical of the traditional TLM structure. This structure additionally allows the measurement of  $R_x$  the cross terminal resistance. The combination of  $R_{||}$  and  $R_x$  enables accurate extractions of  $R_c$ ,  $R_{sh}$ , and  $R_{shm}$ . In actual measurements, the forcing and sensing ports are interchanged so that four sets measurements per elemental test structure can be performed.

Fig. 2 shows a distributed circuit representation of the intrinsic structure of Fig. 1(a). Since Kelvin measurements eliminate the resistance of the extrinsic metal lines, pads and probes, these elements need not be modeled. However, unlike conventional TLMs in which the metal resistance in the

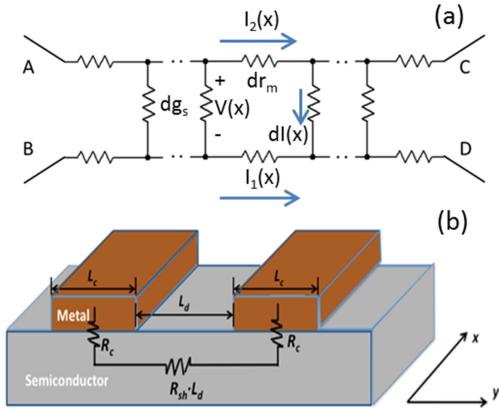


Fig. 2. (a) 2D distributed resistive circuit network of nano-TLM. (b) Cross-sectional view of a nano-TLM with equivalent circuit network.

intrinsic device is not modeled, here it is taken into account. In Fig. 2(a),  $dr_m$  is the elemental longitudinal resistance of the metal line given by:

$$dr_m = \frac{R_{shm}}{L_c} dx \quad (1)$$

$dg_s$  is the elemental metal to metal conductance through the contacts and the semiconductor. Referring to the cross-sectional view of the nano-TLM structure of Fig. 2,  $dg_s$  can be expressed as:

$$dg_s = \frac{1}{R_{TLM}} dx \quad (2)$$

where  $R_{TLM}$  is given as in the conventional TLM [8] :

$$R_{TLM} = 2R_c + R_{sh} \cdot L_d \quad (3)$$

here  $R_c$  is the normalized contact resistance. An analysis of the distributed model of Fig. 2 yields the following expressions for the terminal resistances defined in Fig. 1(b):

$$R_{||} = \frac{R_{TLM}}{L_{Tx}} \csc h \left( \frac{W}{L_{Tx}} \right) \quad (4)$$

$$R_x = \frac{R_{TLM}}{2L_{Tx}} \left[ \csc h \left( \frac{W}{L_{Tx}} \right) + \coth \left( \frac{W}{L_{Tx}} \right) \right] - \frac{R_{shm} W}{2L_c} \quad (5)$$

where  $L_{Tx}$  is the transfer length in the x direction:

$$L_{Tx} = \sqrt{\frac{L_c \cdot R_{TLM}}{2R_{shm}}} \quad (6)$$

As  $R_{shm}$  goes to zero, both (4) and (5) converge to  $R_{TLM}/W$ , which coincides with the classic normalized TLM result.

Fig. 3(a) shows model predictions for  $R_{||}$  and  $R_x$  as a function of contact separation for different contact lengths. In long contacts,  $R_{||}$  and  $R_x$  approach each other. In short contacts, however,  $R_x$  drops below  $R_{||}$  due to the negative term in (5) which captures the voltage drop along the thin metal line.

The model provides guidance to test structure design for accurate extraction of the contact resistance. In general, one would like to have many test structures with identical  $L_c$  and different  $L_d$ . In particular, one needs: (1) a thick metal layer

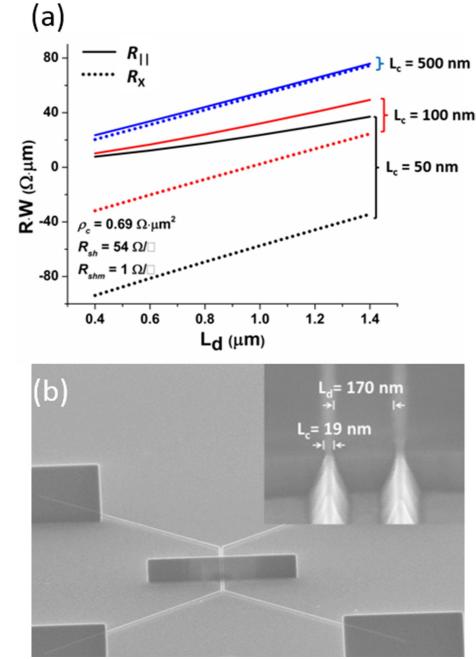


Fig. 3. (a) Modeled terminal resistances versus contact spacing for various contact lengths. (b) SEM image of fabricated Mo/n<sup>+</sup>-InGaAs nano-TLM structure. The inset shows a nano-TLM test structure with 19 nm long contact and 170 nm contact spacing.

to minimize  $R_{shm}$ , (2) small values of  $L_d$  to reduce the impact of  $R_{sh}$ , and (3) a small value of  $W$  to also mitigate the impact of  $R_{shm}$ . The model above suggests that  $W$  should be on the order of  $L_{Tx}$  which in our system is in the range of 1–4  $\mu\text{m}$ . Reasonably small  $W$  also makes the mesa alignment easy. In our devices we have selected  $W = 3 \mu\text{m}$ .

### III. EXPERIMENTS

We have demonstrated the proposed nano-TLM structure in the Mo/n<sup>+</sup>-In<sub>0.53</sub>Ga<sub>0.47</sub>As (20 nm thick, Si-doped,  $N_D = 1 \times 10^{19} \text{ cm}^{-3}$ ) system. We first sputter 60 nm Mo onto the clean fresh semiconductor surface. This metal-first approach has been found to give outstanding contact resistance in InGaAs HEMTs and MOSFETs [10], [11]. The Mo nano-contacts are defined by E-beam lithography followed by RIE. The entire structure is then covered by CVD SiO<sub>2</sub> and the mesas are patterned and isolated by RIE. The depth of mesa etch is 130 nm, which is sufficient for current confinement since only the top 20 nm cap layer is conductive. After pad formation, the structure is annealed at 350 °C in N<sub>2</sub> for 1 min. We have fabricated test structures with contact lengths in the 19 to 450 nm range and contact spacings between 1.4  $\mu\text{m}$  and below. Fig. 3(b) shows a test structure with 19 nm Mo contacts. All the relevant dimensions are measured by scanning electron microscope in the finished devices for accurate analysis.

In each elemental test structure, the four Kelvin measurements mentioned above are performed. The discrepancy in the two measurements of  $R_{||}$  or  $R_x$  is less than 1%. From measurements of test structures with identical nominal  $L_c$  and different  $L_d$  (we call this a “set”), using a nonlinear least-square fitting program, we extract the values of  $R_c$ ,  $R_{sh}$  and  $R_{shm}$  that minimize the error. For accurate extraction,  $R_{||}$

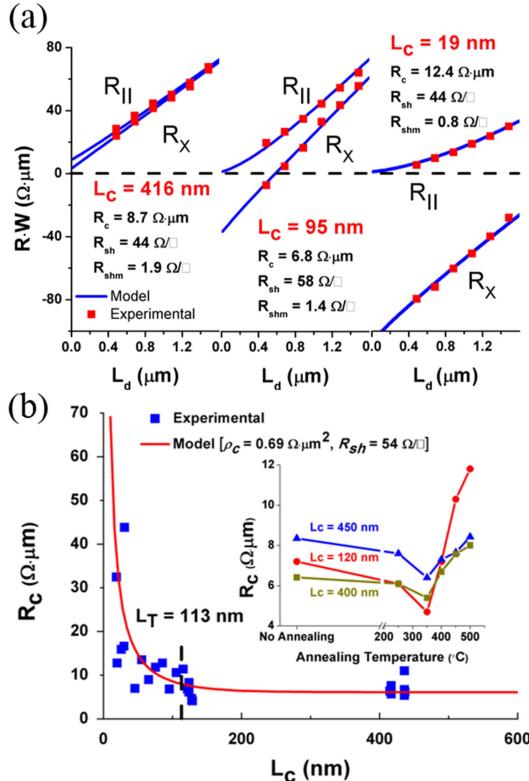


Fig. 4. (a)  $R_{\parallel}$  and  $R_X$  of nano-TLMs with average contact lengths of 416, 95, and 19 nm. (b) Contact resistance versus contact length for  $\text{Mo}/\text{n}^+ \text{-In}_{0.53}\text{Ga}_{0.47}\text{As}$  contacts, with an average contact resistivity of  $0.69 \pm 0.3 \Omega \cdot \mu\text{m}^2$ . The inset shows the thermal stability of the Mo contacts.

and  $R_X$  are fitted simultaneously. Examples from three sets of test structures with  $L_c = 416$ , 95, and 19 nm are shown in Fig. 4(a). Even for relatively large 95 nm contacts, there is a remarkable split between the two measurement schemes,  $R_{\parallel}$  and  $R_X$ , and they split further for smaller contacts. Moreover, while  $R_{\parallel}$  is always positive,  $R_X$  can become negative for small contacts, consistent with the prediction of the model in Fig. 3(a).

Since we expect non-uniformities in test structure fabrication, we allow  $R_{sh}$  and  $R_{shm}$  to change from test structure set to set. The obtained values of  $R_{sh}$  and  $R_{shm}$  have average values of  $54 \Omega/\square$  and  $1.2 \Omega/\square$ , and standard deviations of  $4.3 \Omega/\square$  and  $0.2 \Omega/\square$ , respectively. To verify the accuracy of our extractions, we also obtained  $R_{sh}$  and  $R_{shm}$  from conventional TLMs and Van der Pauw structures elsewhere on the sample. The average discrepancy of  $R_{sh}$  and  $R_{shm}$  are 10% and 16% across the sample.

Fig. 4(b) shows the dependence of  $\text{Mo}/\text{n}^+ \text{-InGaAs}$  contact resistance on contact length obtained in 30 sets of test structures, a collective of 180 nano-TLM devices. For relatively long contacts, an extremely small average  $R_c = 6.6 \pm 1.6 \Omega \cdot \mu\text{m}$  is achieved.  $R_c$  shoots up as the contact length becomes smaller than the transfer length in the y direction:

$$L_{Ty} \approx \sqrt{\rho_c / R_{sh}} \quad (7)$$

which is about 113 nm for this contact system, assuming that  $R_{sh} \gg R_{shm}$ . We have used the standard TLM model as in [8] to extract the contact resistivity, and we find

$\rho_c = 0.69 \pm 0.3 \Omega \cdot \mu\text{m}^2$  for all contact lengths. Fig. 4(b) shows that the simple model using the extracted  $\rho_c$  and  $R_{sh}$  provides a good match to the entire data set. In the extraction of  $\rho_c$ , we assume the semiconductor sheet resistance remains the same under the metal contact [12]. This value of  $\rho_c$  supersedes the one erroneously reported in [13] due to a flawed extraction. The inset of Fig. 4(b) also shows sequential measurements of  $R_c$  after  $\text{N}_2$  annealing for 1 min at increasing temperature. These results suggest that Mo/InGaAs contacts are thermally stable up to at least 400 °C, in good agreement with [7], [11].

#### IV. CONCLUSION

We have proposed a new two-port TLM test structure to characterize nano-scale contacts. We demonstrate this concept in  $\text{Mo}/\text{n}^+ \text{-InGaAs}$  contacts with contact lengths in the 19–450 nm range. With the help of a 2D distributed resistive circuit network, we have extracted the contact resistance of contacts as short as 19 nm. Our entire data set is well described by a contact resistivity of  $\rho_c = 0.69 \pm 0.3 \Omega \cdot \mu\text{m}^2$  which is close to the requirements for future nanometer-scale CMOS.

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