

# Breakdown Voltage Enhancement from Channel Quantization in InAlAs/ $n^+$ -InGaAs HFET's

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**Abstract**—In<sub>0.52</sub>Al<sub>0.48</sub>As/ $n^+$ -In<sub>0.53</sub>Ga<sub>0.47</sub>As HFET's have been fabricated with different channel thicknesses. We show that by reducing the channel thickness from 350 to 100 Å, the reverse gate breakdown voltage improves from 9 to 19 V. We partially attribute this to the increased effective bandgap that results from energy quantization in the channel. This bandgap enhancement is directly confirmed by photoluminescence (PL) measurements on the same heterostructures. Channel quantization emerges as a promising approach for exploiting the excellent transport properties of InGaAs with high InAs mole fraction. The principle behind our work should be applicable to other narrow-gap semiconductors.

In<sub>0.52</sub>Al<sub>0.48</sub>As/In<sub>0.53</sub>Ga<sub>0.47</sub>As heterostructure field-effect transistors (HFET's) on InP are of great interest for long-wavelength optical and ultrahigh-frequency microwave telecommunications applications. The advantages of this material system are many. In<sub>0.53</sub>Ga<sub>0.47</sub>As has a higher peak electron velocity and a higher room-temperature mobility than GaAs. In addition, it has a larger  $\Gamma$ - $L$  separation and a lower effective mass [1].

Enriching the InAs mole fraction of In<sub>0.53</sub>Ga<sub>0.47</sub>As results in HFET's with superior transport properties. However, this comes at the cost of a severely reduced breakdown voltage  $V_B$ , presumably through the decrease in the energy gap  $E_g$  [2]. A method of increasing the effective energy gap in the channel is to introduce energy quantization by reducing the channel thickness to dimensions comparable to the electron wavelength (Fig. 1) [3]. In fact, it has been shown in In<sub>0.53</sub>Ga<sub>0.47</sub>As/In<sub>0.52</sub>Al<sub>0.48</sub>As quantum wells that the photoluminescence emission wavelength decreases [4], [5] with a reduction in well thickness.

In this work, we exploit this effect to enhance the breakdown voltage of In<sub>0.52</sub>Al<sub>0.48</sub>As/ $n^+$ -In<sub>0.53</sub>Ga<sub>0.47</sub>As HFET's on InP. We have doubled  $V_B$  by shrinking the In<sub>0.53</sub>Ga<sub>0.47</sub>As channel thickness from 350 to 100 Å, keeping other physical

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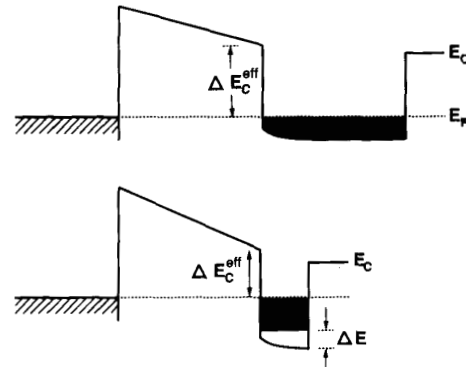


Fig. 1. Schematic conduction band diagrams in equilibrium of InAlAs/ $n^+$ -InGaAs HFET's for thick and thin channels, showing the increase in effective energy gap in the thinner channel.

parameters constant. The principle behind our work should allow one to better exploit the excellent transport properties of InAs-rich InGaAs [2] and other promising narrow-gap semiconductors like InAs and InSb [6], [7].

A cross section of the device structure is shown in Fig. 2. In essence, this device is a doped-channel HFET with an undoped pseudoinulator [8]. The device structure consists (from bottom to top) of a 1000-Å undoped In<sub>0.52</sub>Al<sub>0.48</sub>As buffer layer, a channel consisting of an undoped In<sub>0.53</sub>Ga<sub>0.47</sub>As layer and a 100-Å heavily Si-doped (nominal  $N_D = 4 \times 10^{18} \text{ cm}^{-3}$ ) In<sub>0.53</sub>Ga<sub>0.47</sub>As layer, a 300-Å undoped In<sub>0.52</sub>Al<sub>0.48</sub>As gate insulator layer, and an undoped 50-Å In<sub>0.53</sub>Ga<sub>0.47</sub>As cap. The wafers were grown on Si-InP by MBE in M.I.T.'s Riber 2300 system. In an effort to keep the channel charge constant, the thickness of its undoped portion was varied while the thickness of its heavily portion was kept constant at 100 Å. Four wafers were subsequently grown with subchannel thicknesses of 250, 100, 50, and 0 Å, i.e., total channel thicknesses of 350, 200, 150, and 100 Å. Devices were fabricated with nominal gate lengths of 1  $\mu\text{m}$  and widths of 30  $\mu\text{m}$ . Fabrication is similar to that used in [8].

Our baseline device, 200-Å channel thickness, had a peak transconductance  $g_{m(\text{peak})}$  of 202 mS/mm and a maximum drain current  $I_{d(\text{max})}$  of 312 mA/mm. The output conductance  $g_d$  was 5.73 mS/mm, resulting in a voltage gain  $A_V$  of 35.

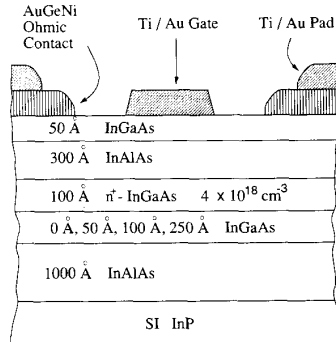


Fig. 2. Cross section of grown device structures.

These are average values over five devices.  $g_{m(\text{peak})}$  and  $I_{d(\text{max})}$  were measured at  $V_{ds} = 4$  V, and  $g_d$  at  $V_{ds} = 4$  V and  $V_{gs} = 0$  V. The contact and channel sheet resistances, measured by TLM, are  $0.37 \Omega \cdot \text{mm}$  and  $625 \Omega/\square$ , respectively.

The reverse gate breakdown voltage  $V_B$  was measured with the source and drain grounded, and was defined at a reverse gate current of  $500 \mu\text{A}$ , corresponding to about 5% of the peak drain current of our baseline device [2]. Detailed device results are presented elsewhere [9]. Here we focus on our main result: the increased breakdown voltage in devices with thinner channels and the experimental confirmation of the energy quantization therein.

Fig. 3 shows typical reverse gate  $I$ - $V$  characteristics of HFET's as a function of channel thickness. As shown,  $V_B$  increases gradually from 9 V at a channel thickness of  $350 \text{ \AA}$  to  $10.6$  V at  $200 \text{ \AA}$ ,  $11.9$  V at  $150 \text{ \AA}$ , and to  $19.1$  V at  $100 \text{ \AA}$ . Average  $V_B$  measurements over several devices are within 1 V of the typical values shown in Fig. 3.  $V_B$  for the  $200\text{-\AA}$  channel is also similar to what we have previously measured in identical devices grown and processed separately [2]. The drastic improvement in breakdown voltage of our quantized-channel HFET's is a significant merit for high-power applications. This is particularly so in this material system because typical InAlAs/InGaAs MODFET breakdown voltages are on the order of 5 V [10], [11].

In order to verify the bandgap enhancement in the channel as a result of carrier quantization, we have carried out photoluminescence (PL) measurements on unprocessed portions of the device samples at 77 K. The results are shown in Fig. 4. The energy of the peak PL intensity increases from  $0.83$  eV for the  $350\text{-\AA}$  channel to  $0.86$  eV for the  $200\text{-\AA}$  channel,  $0.88$  eV for the  $150\text{-\AA}$  channel, and  $0.92$  eV for the  $100\text{-\AA}$ -channel. The literature reports a temperature-independent PL energy shift of about  $60$  meV over the  $\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$  bulk value for  $100\text{-\AA}$ -thick  $\text{In}_{0.52}\text{Al}_{0.48}\text{As}/\text{In}_{0.53}\text{Ga}_{0.47}\text{As}$  quantum wells [4], [5]. This value of  $60$  meV is also consistent with simple calculations for a finite square well of thickness  $100 \text{ \AA}$ . Our slightly larger shift of  $90$  meV in the  $100\text{-\AA}$  channel might be due to depletion at the top and possibly bottom interfaces, causing a reduced effective well thickness. Additionally, band bending in the insulator results in an effective stronger potential than a square well (Fig. 1).

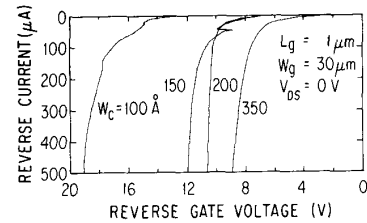


Fig. 3. Breakdown voltage  $V_B$  for typical HFET's with channel thicknesses of 100, 150, 200, and  $350 \text{ \AA}$ .

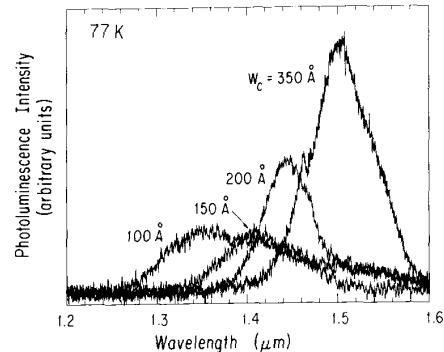


Fig. 4. PL spectra of the device heterostructures, showing an increase in PL energy with decreasing channel thickness.

Both effects will tend to enhance the strength of carrier quantization over a simple quantum well.

Electron quantization, as schematically shown in Fig. 1, also implies a reduction of the effective conduction-band discontinuity between channel and metal gate. This, in return, is expected to result in enhanced forward gate leakage current. We have experimentally found this to be the case at 300 and 77 K [9], providing us in this manner with an independent confirmation of the presence of quantization in the channel.

Unfortunately, we have found that a reduction in the channel thickness results in the degradation of transconductance and peak drain current [9].  $g_{m(\text{peak})}$  decreases from  $262$  to  $138$  mS/mm, and  $I_{d(\text{max})}$  from  $451$  to  $208$  mA/mm in going from a channel thickness of  $350$  to  $100 \text{ \AA}$ . However, the output conductance  $g_d$  improves due to the enhanced channel aspect ratio, decreasing from  $10.8$  to  $1.84$  mS/mm. This results in an enhancement in the voltage gain  $A_V$  from  $24$  to  $75$ . The degradation in  $g_{m(\text{peak})}$  and  $I_{d(\text{max})}$  results from an increased source resistance  $R_s$ , a reduced channel sheet charge concentration  $n_s$ , and degraded mobility  $\mu$ . From  $350$  to  $100 \text{ \AA}$ ,  $R_s$  increases from  $1.4$  to  $2.7 \Omega \cdot \text{mm}$ ,  $n_s$  drops from  $2.38 \times 10^{12}$  to  $1.77 \times 10^{12} \text{ cm}^{-2}$ , and  $\mu$  decreases from  $4318$  to  $3591 \text{ cm}^2/\text{V} \cdot \text{s}$ .  $n_s$  and  $\mu$  were measured by the Hall effect. The acknowledged poor quality of the reverse InGaAs/InAlAs interface at the back of the channel can be held responsible for the mobility reduction [12], [13]. As the undoped subchannel is thinned down, the reverse interface has a larger impact on carrier mobility since the channel electrons travel closer to the reverse InAlAs/InGaAs interface. There are, however, techniques that could mitigate this

degradation: superlattice buffers to improve mobility [13] and the migration-enhanced epitaxy (MEE) growth technique to reduce interface roughness [14].

Thinning down the subchannel results in a reduction of sheet carrier concentration, possibly from backside depletion. A reduced channel doping can also result in an improvement of breakdown voltage. In a separate experiment on similar devices (with 200-Å channel thickness), we examined the impact of channel doping on  $V_B$ . This experiment indicated that a reduction in sheet charge concentration from  $2.38 \times 10^{12}$  to  $1.77 \times 10^{12} \text{ cm}^{-2}$  should result in an improvement of  $V_B$  by 5 V. Our experimental observation of a 10-V improvement in going from a 350- to 100-Å channel is evidence that quantization is instrumental in drastically enhancing the breakdown characteristics of our HFET's.

In conclusion,  $\text{In}_{0.52}\text{Al}_{0.48}\text{As}/\text{n}^+\text{-In}_{0.53}\text{Ga}_{0.47}\text{As}$  HFET's have been fabricated with channel thicknesses of 100, 150, 200, and 350 Å. For devices with channel thicknesses of 100 Å, the breakdown voltage improved twofold over the 350-Å devices. This is postulated to partially arise from an enlargement of the effective energy gap caused by energy quantization introduced from electron confinement, as observed by PL.

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